

## **X-ray constraints on late decaying dark matter majorons - or other soft X-ray emitting candidates**

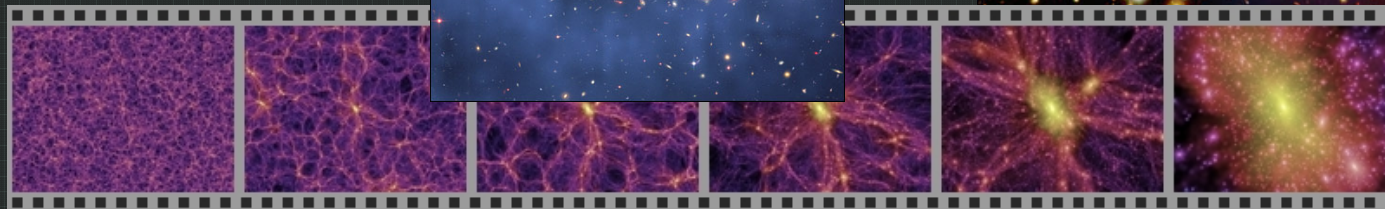
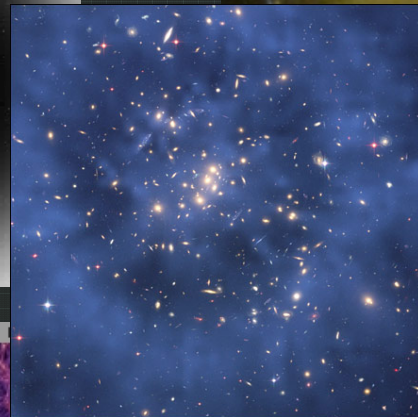
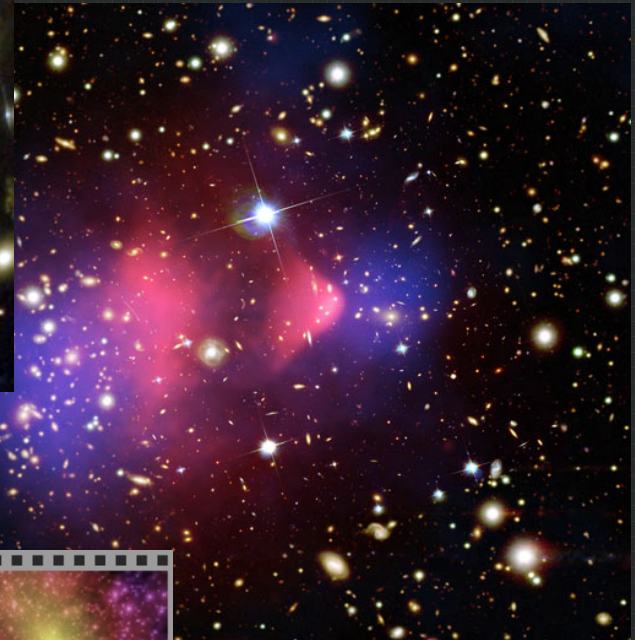
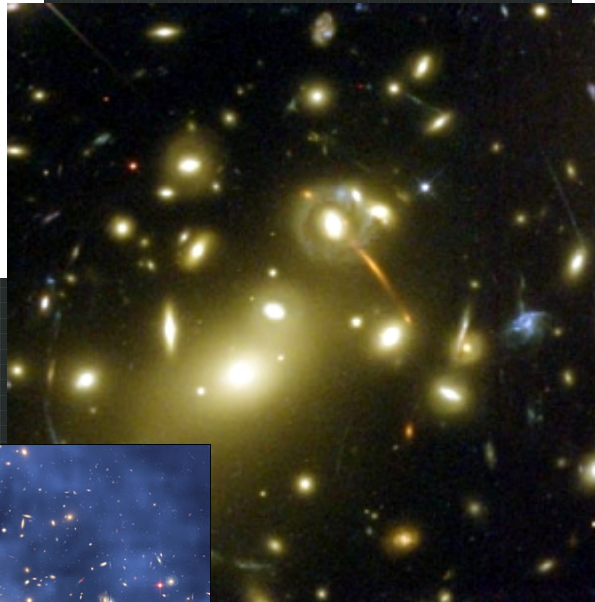
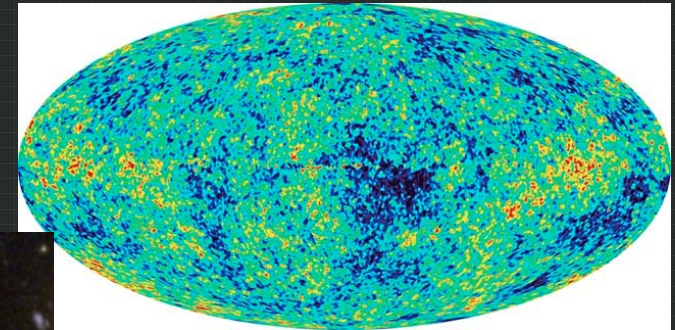
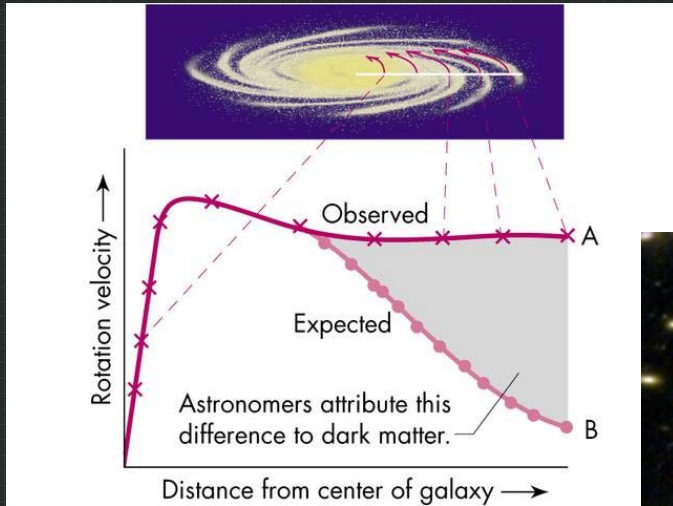
4th Patras Workshop on Axions, WIMPs and WISPs, DESY, June 18-21 2008

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arXiv:astro-ph/0805.2372, accepted for JCAP:

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# Something is missing

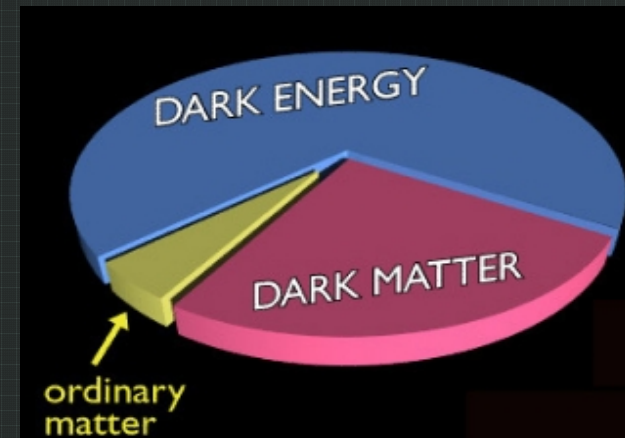


General properties of a dark matter candidate:

- Particle behavior
- Massive (gravitational effect)
- Not too much interacting
- Long lifetime (if thermal relic)

No Standard Model candidate!

Cosmology -> only 4% Standard Model



Experimental evidence for problems with the Standard Model e.g.:

- Neutrino flavour mixing
- Neutrino masses
- Baryon asymmetry
- Dark matter
- Dark energy

Neutrinos massless in the Standard Model  
BUT experimental evidence for  
masses

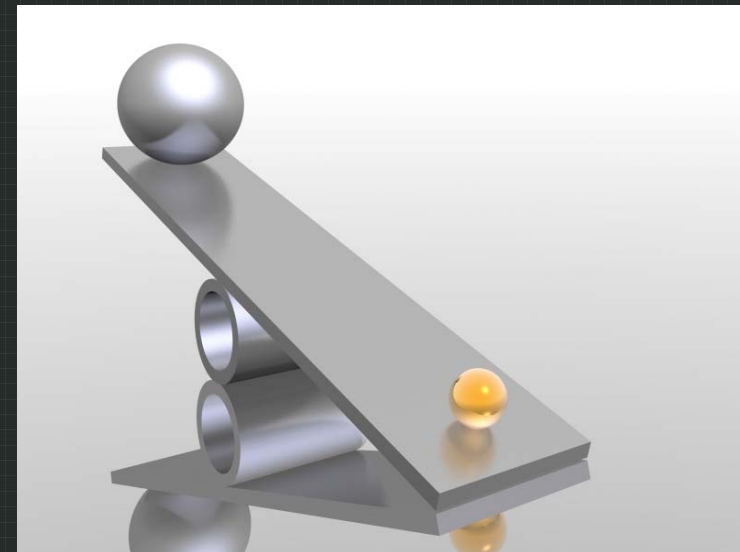
Seesaw mechanism popular to generate  
masses

Two types of neutrino mass terms:

Dirac particles -> Lepton number  
conserving

Majorana particles -> Lepton number  
violating

Global lepton symmetry, neutrino masses  
acquired by spontaneous violation



Neutrino masses can arise from the spontaneously breaking of un-gauged lepton number

Pseudoscalar massless Nambu-Goldstone gauge boson: The majoron

(Chikashige, Mohapatra & Peccei 1981, Schechter & Valle 1982)

Majoron acquire mass from non-perturbative gravitational effects that explicitly break global symmetries

(Akhmedov et al. 1993)

Dark matter: Non Standard Model, massive, little interacting

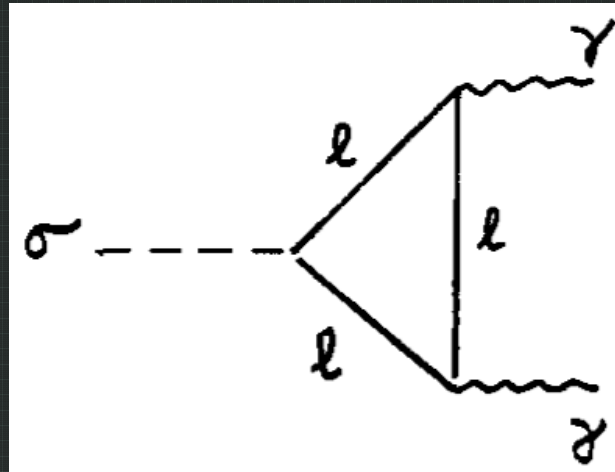
(Berezinsky & Valle 1993)

Dominating decay into neutrinos

Can decay to photons - loop suppressed

(Berezinsky & Valle 1993)

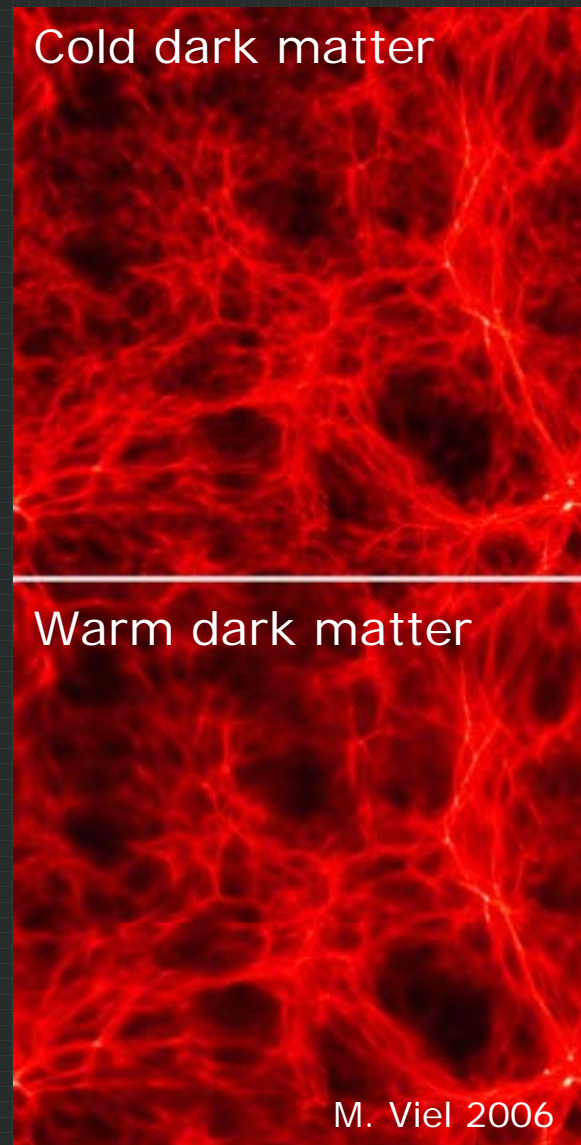
$$J \rightarrow \gamma \gamma$$
$$E_\gamma = m_{J/2}$$



IF produced thermally and very early:  
keV mass provides present day dark matter  
density

(Lattanzi & Valle, 2006)

Warm dark matter



Late decay of majorons change gravitational potential

Late integrated Sachs-Wolfe effect causes excess of power at small multipoles (in CMB)

(Lattanzi & Valle 2007)

Constraints:

$$0.11 \text{ keV} < \beta m_J < 0.18 \text{ keV}$$

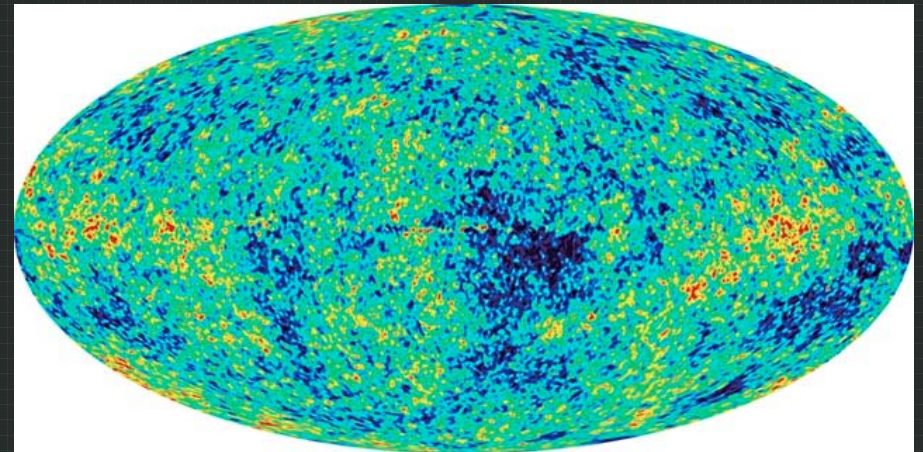
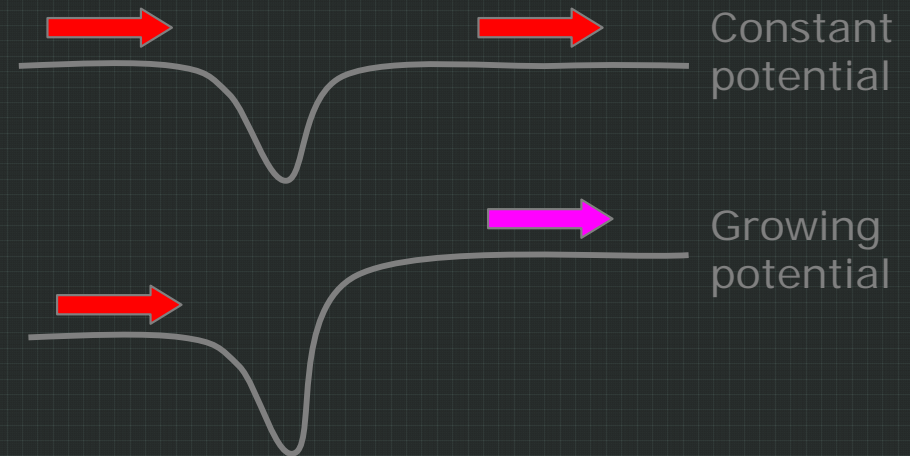
$$t_J > 250 \text{ Gyr}$$

$$\Gamma_{J\nu\nu} < 1.3 \cdot 10^{-19} \text{ s}^{-1}$$

(Lattanzi & Valle 2007)

Assumptions:

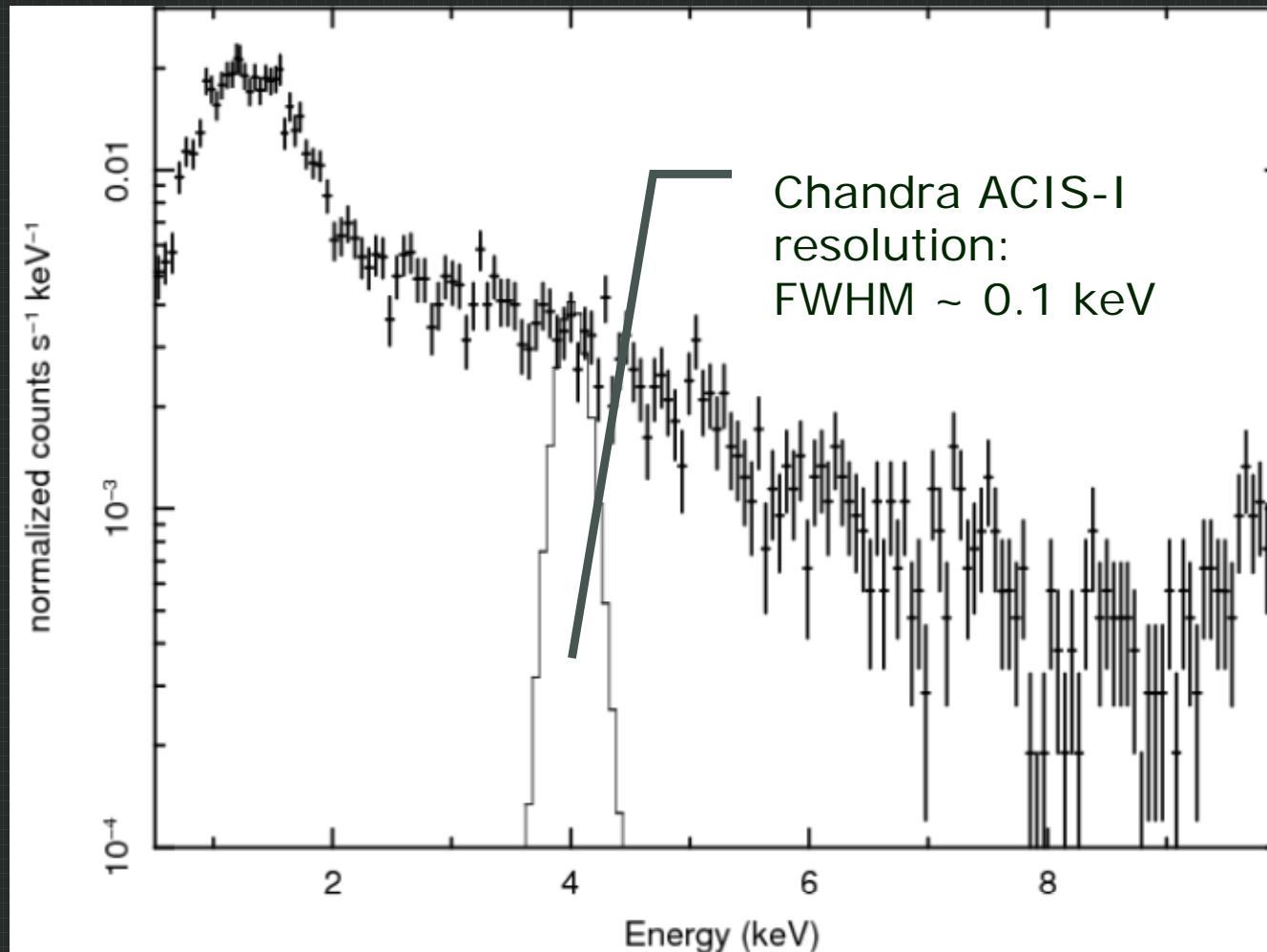
- Thermally produced and in equilibrium in early Universe
- Decouple while all SM degrees of freedom are excited





Dark matter decaying into photons,  $E_\gamma = m_{J/2}$

Mono-energetic emission -> Gaussian because of instrumental resolution



$$F \propto \frac{M\Gamma}{D_L^2}$$

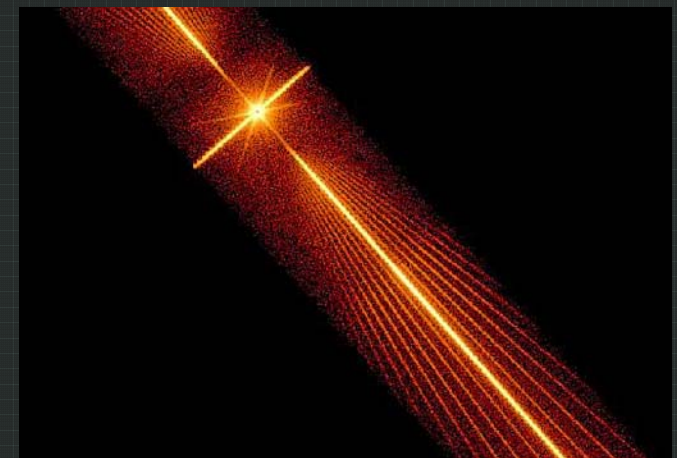
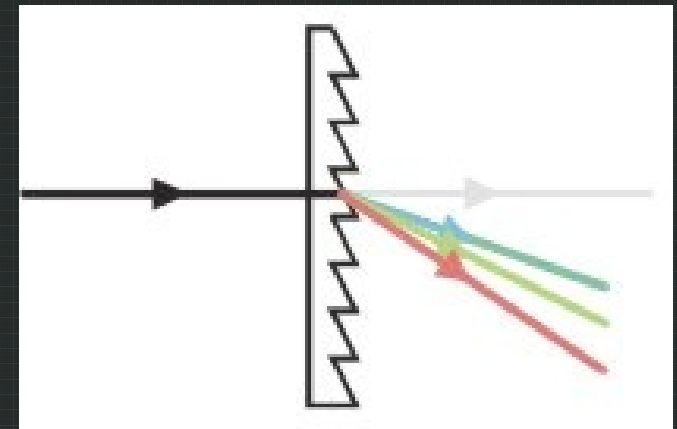
X-ray absorbed by atmosphere -> space borne instruments

Chandra and XMM imaging spectrographs only sensitive down to 0.3 keV

BUT Chandra carries an imaging instrument sensitive down to 0.07 keV,  
which unfortunately can't do spectra...

Unless combined with a grating:

- Separate wavelengths -> resolution of 0.005 keV for point sources
- Extended sources reduces resolution (line broadening)
- All information on origin of photon is lost
- Requires bright sources

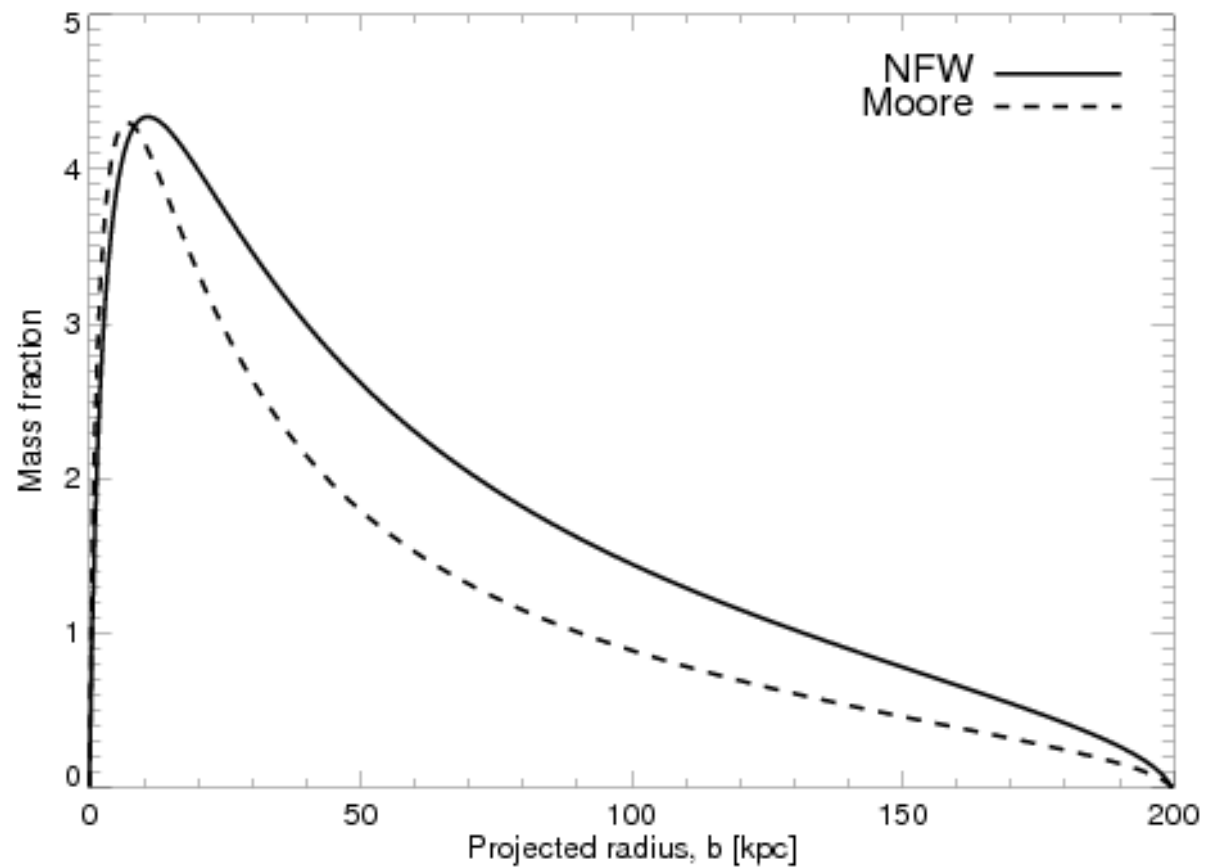
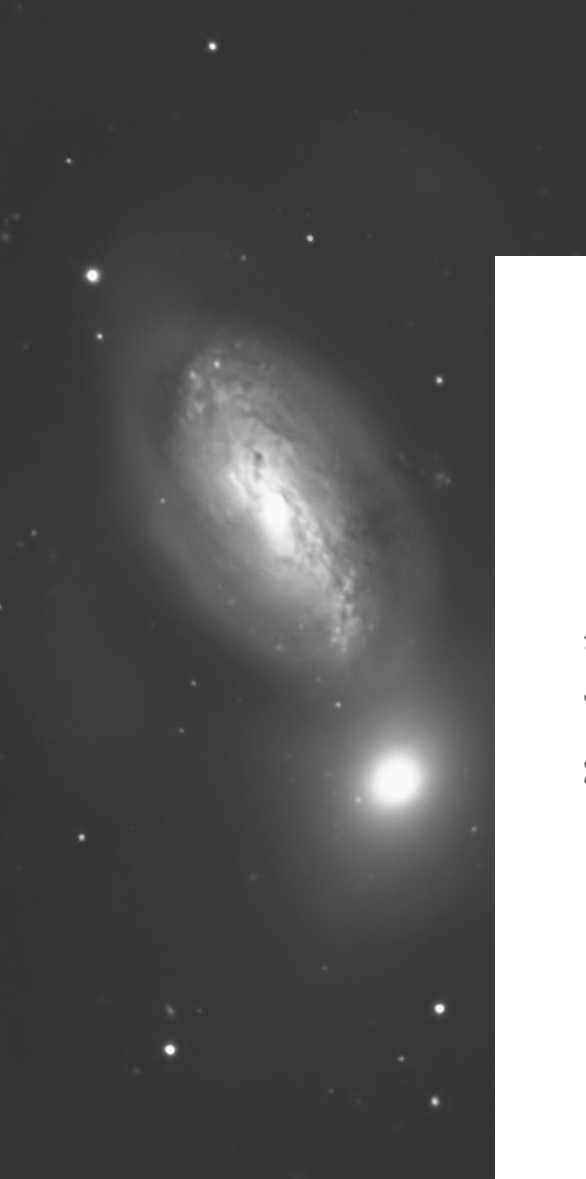


# Observations of NGC3227

NGC3227 is Seyfert 1 galaxy at  $z=0.004$

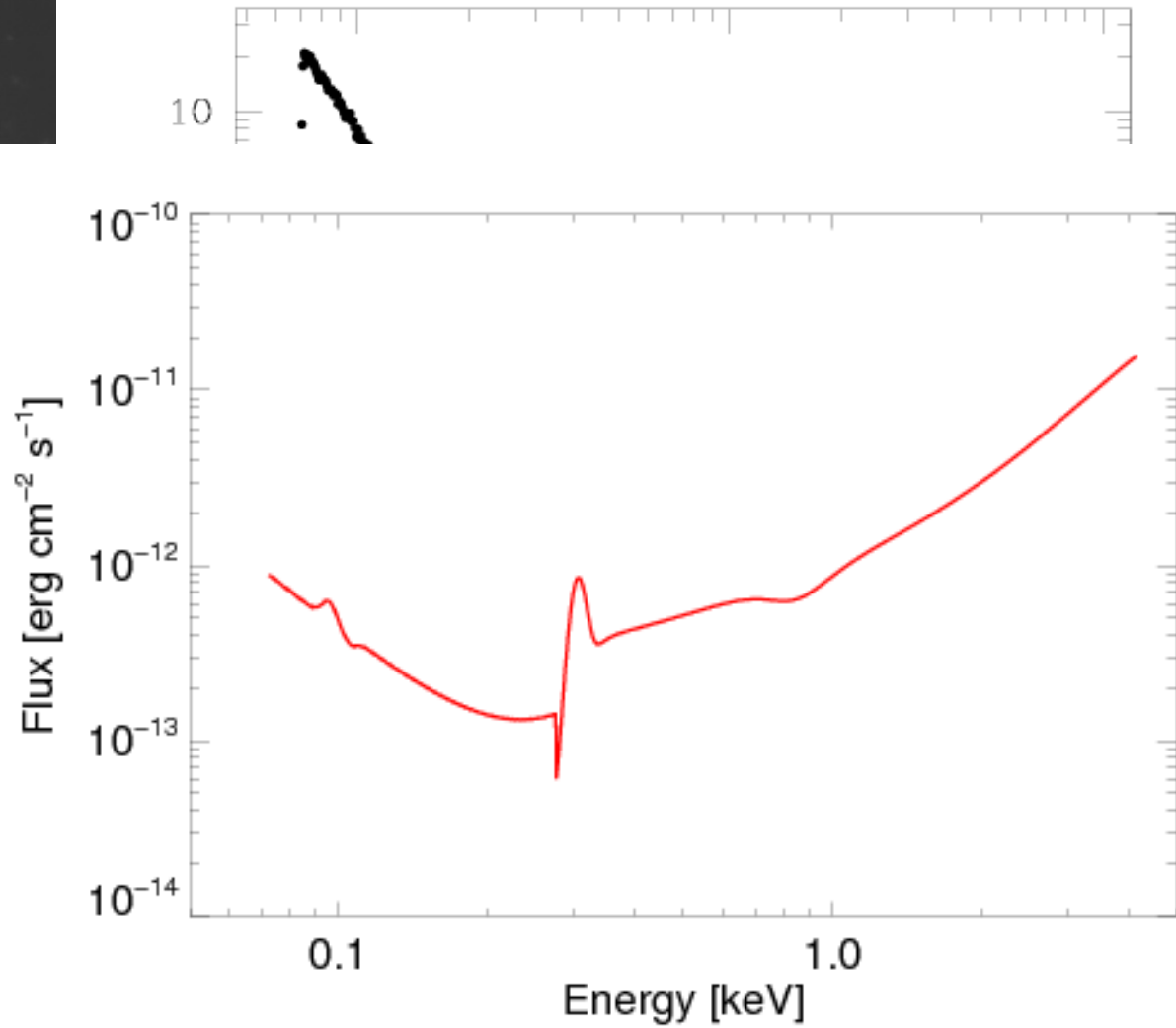
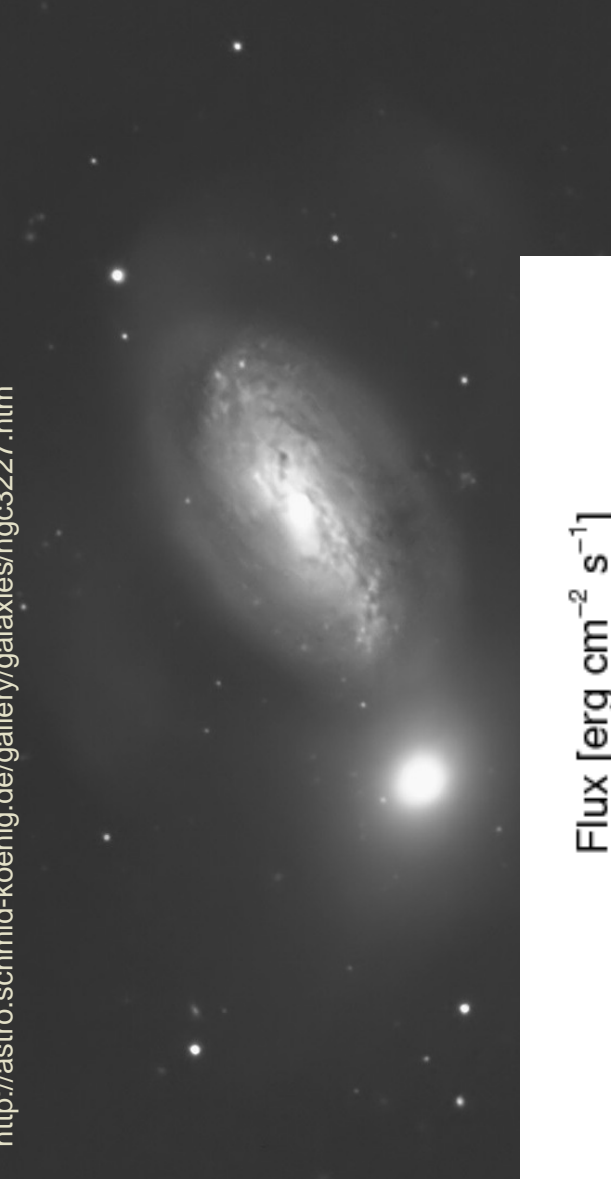
Typical mass of  $\sim 10^{11}-10^{12} M_{\text{sun}}$ , 1/10 in field of view, approx. 80% dark matter

<http://astro.schmid-koenig.de/gallery/galaxies/ngc3227.htm>

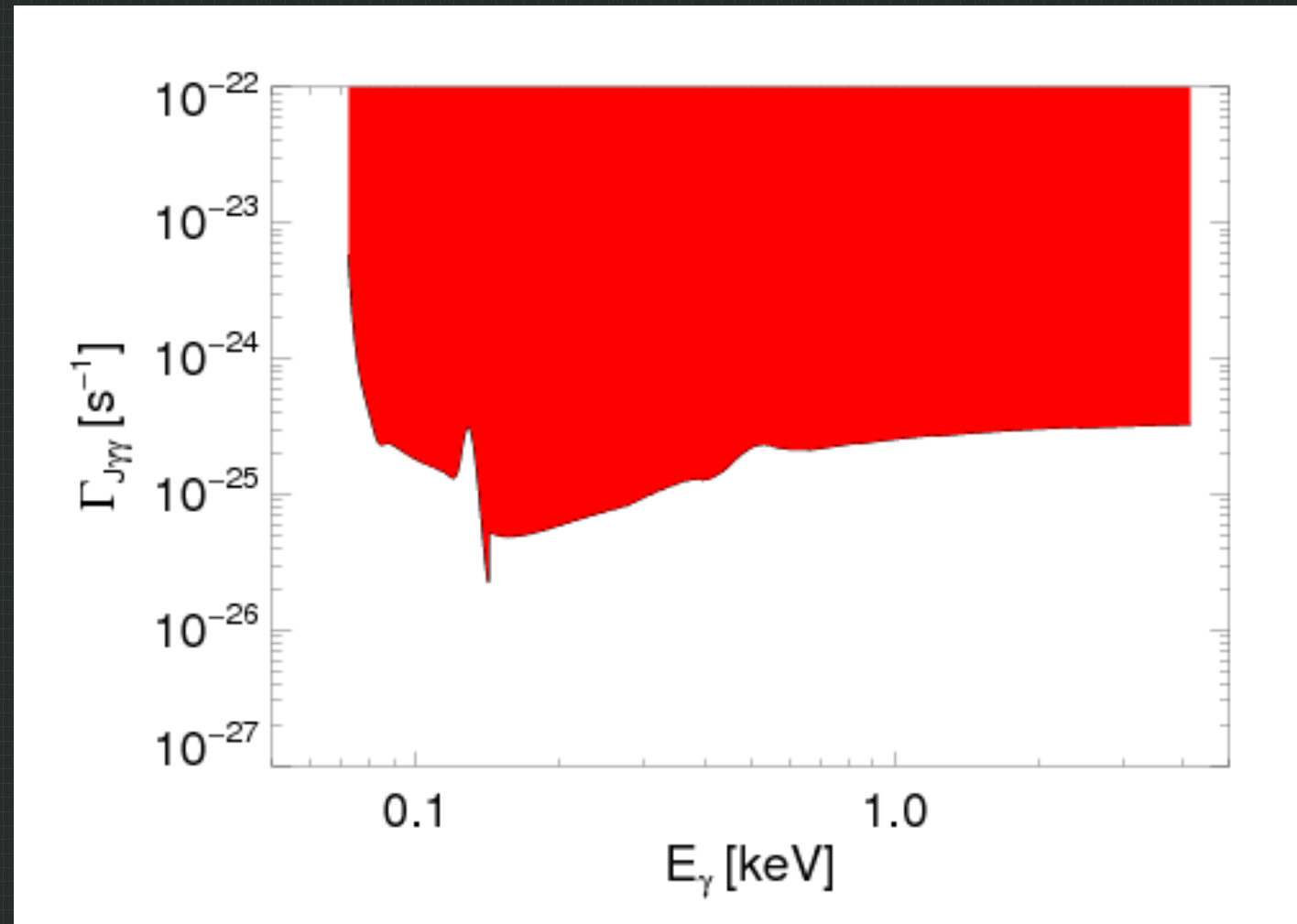


# Observed flux

<http://astro.schmid-koenig.de/gallery/galaxies/ngc3227.htm>



$\Gamma_{J\gamma} < 10^{-24} \text{ s}^{-1}$  consistent with  $\Gamma_{J\gamma} \ll \Gamma_{J\nu\nu} < 1.3 \cdot 10^{-19} \text{ s}^{-1}$



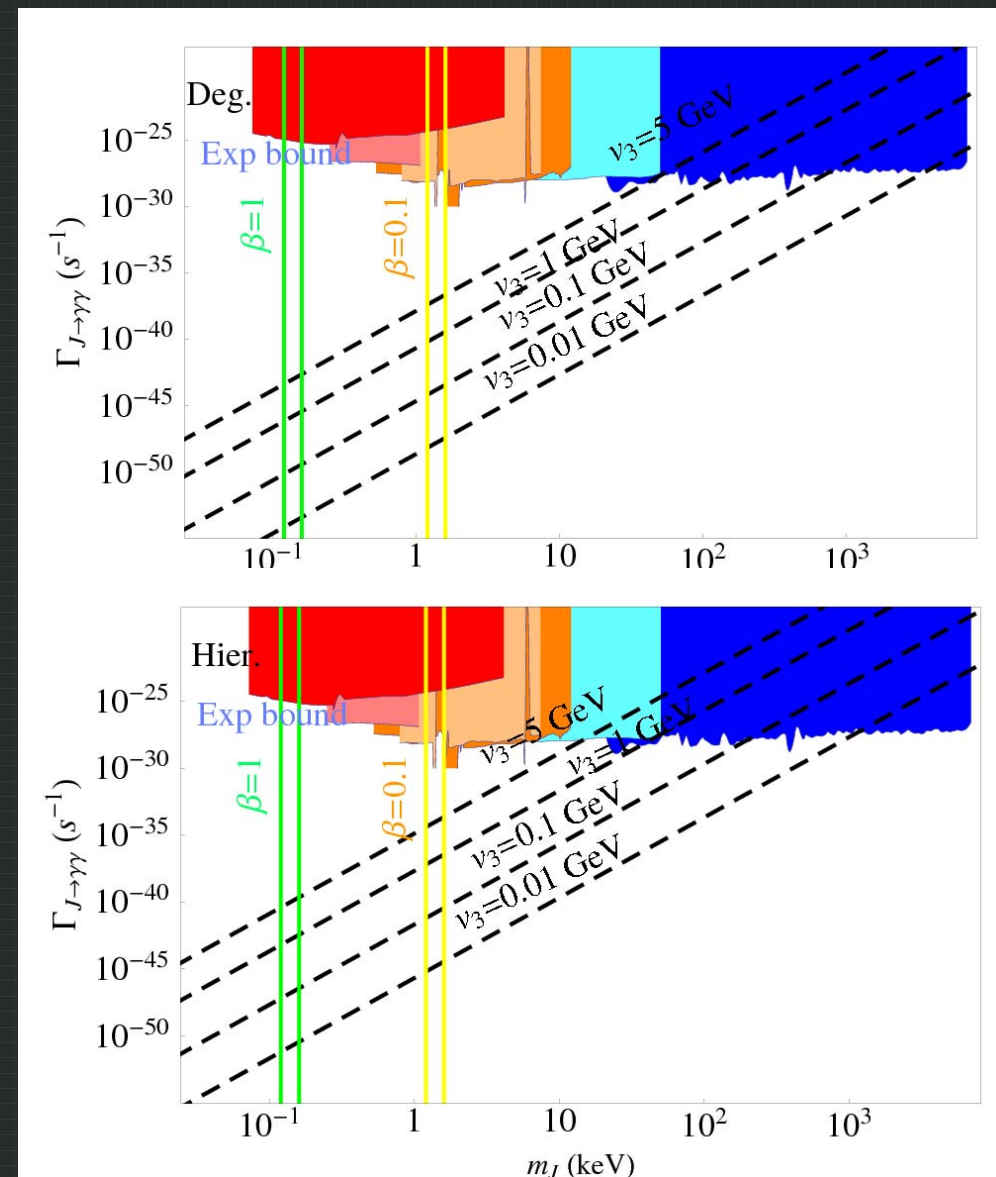
# Relaxing assumptions

CMB constraints:

$$0.11 \text{ keV} < \beta m_J < 0.18 \text{ keV}$$

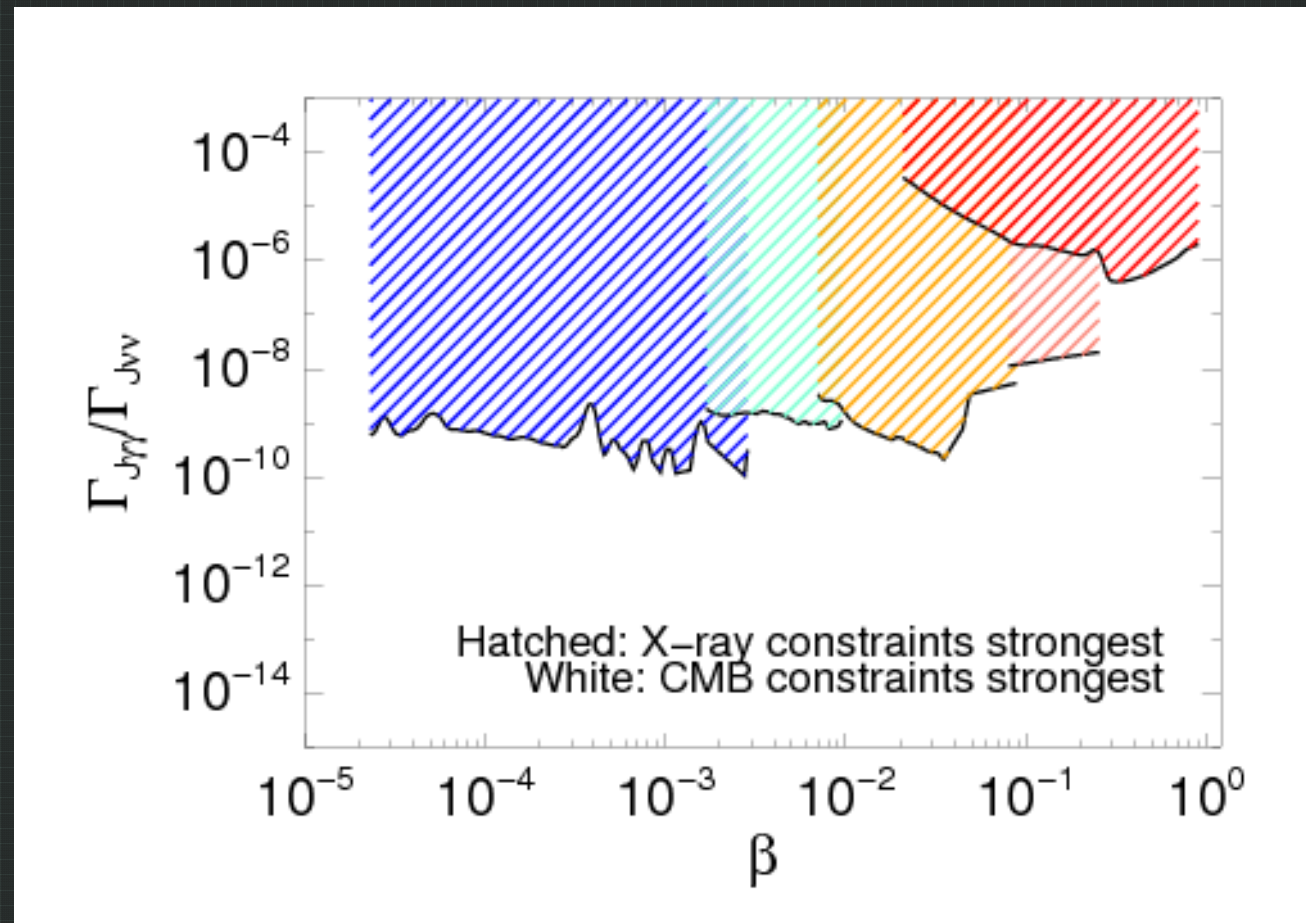
Relaxing assumptions  
about thermal  
equilibrium and early  
decoupling switches the  
relevant mass scale  
(change  $\beta$ )

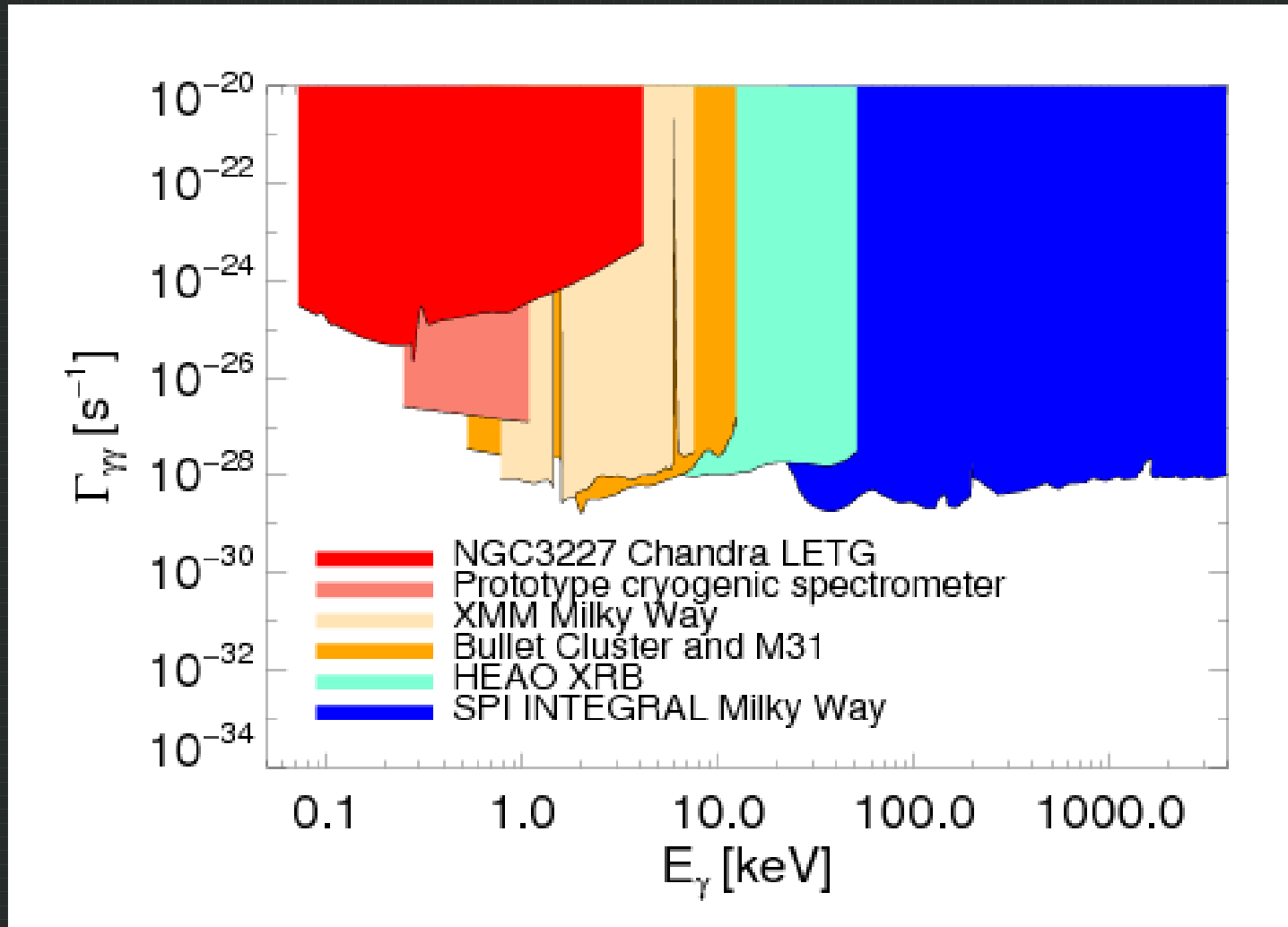
Boyarsky et al. (2007) *Astropart. Phys.* 28, 303 (salmon), Boyarsky et al. (2007) *Astron. Astrophys.* 471, 51 (sand), Boyarsky et al. (2008) *Astrophys. J.* 673, 752 & Boyarsky et al. (2007), arXiv:0709.2301 (orange), Boyarsky et al. (2006) *MNRAS* 370, 213 & Boyarsky et al. (2006) *Phys. Rev. Lett.* 97, 261302 (aquamarine), Boyarsky et al. (2007) arXiv:0710.4922 (blue)



# Which constraint is stronger

Depends on branching ratio and  $\beta$





Applies also to  
your favorite  
candidate!

Boyarsky et al. (2007) *Astropart. Phys.* 28, 303 (salmon), Boyarsky et al. (2007) *Astron. Astrophys.* 471, 51 (sand), Boyarsky et al. (2008) *Astrophys. J.* 673, 752 & Boyarsky et al. (2007), arXiv:0709.2301 (orange), Boyarsky et al. (2006) *MNRAS* 370, 213 & Boyarsky et al. (2006) *Phys. Rev. Lett.* 97, 261302 (aquamarine), Boyarsky et al. (2007) arXiv:0710.4922 (blue)



Majorons from lepton number violation might be massive -> DM candidate

Late integrated Sachs-Wolfe effect points towards  $\sim 0.15$  keV mass

Observationally constrained decayrate - unfortunately far from sensitivity needed to constrain models

Wide range of observations - applies to all dark matter candidates with a two-body radiative decay

